

Putting a Spin on Language: A Quantum Interpretation of Unary Connectives for Linguistic Applications

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Extended versions of the Lambek Calculus currently used in computational linguistics rely on unary modalities to allow for the controlled application of structural rules affecting word order and phrase structure. These controlled structural operations give rise to derivational ambiguities that are missed by the original Lambek Calculus or its pregroup simplification. Proposals for compositional interpretation of extended Lambek Calculus in the compact closed category of \mathbf{FVect} and linear maps have been made, but in these proposals the syntax-semantics mapping ignores the control modalities, effectively restricting their role to the syntax. Our aim is to turn the modalities into first-class citizens of the vectorial interpretation. Building on the density matrix semantics of (Correia et al, 2019), we extend the interpretation of the type system with an extra spin density matrix space. The interpretation of proofs then results in ambiguous derivations being tensored with orthogonal spin states. Our method introduces a way of simultaneously representing co-existing interpretations of ambiguous utterances, and provides a uniform framework for the integration of lexical and derivational ambiguity.

1 Introduction

A cornerstone of formal semantics is Montague’s [12] compositionality theory. Compositional interpretation, in this view, is a homomorphism, a structure-preserving map that sends types and derivations of a syntactic source logic to the corresponding semantic spaces and operations thereon. In the DisCoCat framework [5] compositionality takes a surprising new turn. Montague’s abstract mathematical view on the syntax-semantics interface is kept, but the non-committed view on *lexical* meaning that one finds in formal semantics is replaced by a data-driven, distributional modelling, with finite dimensional vector spaces and linear maps as the target for the interpretation function. More recently density matrices and completely positive maps have been used to treat lexical ambiguity [17], word and sentence entailment [2, 18] and meaning updating [3].

Our goal in this paper is to apply the DisCoCat methodology to an extended version of the Lambek calculus where structural rules affecting word order and/or phrase structure are no longer freely available, but have to be explicitly licensed by unary control modalities, [9, 13]. In particular, we adjust the interpretation homomorphism to assign appropriate semantic spaces to the modally extended type language, and show what their effect is on the derivational semantics. We choose to use density matrices as our interpretation spaces and show that, besides allowing for an integration of our model with other forms of ambiguity at the lexical level, it is key to preserve information about the ambiguity at phrase level.

The paper is structured as follows. In section 2 we recall the natural deduction rules of the simply typed Lambek Calculus, with the associated lambda terms under the proofs-as-programs interpretation. We extend the language with a residuated pair of unary modalities \diamond, \square and show how these can be used to control structural reasoning, in particular reordering (commutativity). As an illustration, we show how the extended type logic allows us to capture derivational ambiguities that arise in Dutch relative clause constructions. In section 3 we set up the mapping from syntactic types to semantic spaces, adding an extra spin space to the previously used density matrix spaces. We

motivate the introduction of this extra space and relate the interpretation of the connectives in these spaces to the measurement and evolution postulates of quantum mechanics. In section 4 we show how the interpretation of the logical and structural inference rules of our extended type logic accommodates the spin space. In section 6 we return to our example of derivational ambiguity and show how orthogonal spin states keep track of co-existing interpretations.

2 Extended Lambek Calculus

By \mathbf{NL}_\diamond we designate the (non-associative, non-commutative, non-unital) pure residuation logic of [11], extended with a pair of unary type-forming operators \diamond, \square , also forming a residuated pair. Formulas are built over a set of atomic types \mathcal{A} (here s, np, n for sentences, noun phrases and common nouns respectively) by means of a binary product \bullet with its left and right residuals $/, \backslash$, and a unary \diamond with its residual \square :

$$\mathcal{F} ::= \mathcal{A} \mid \square \mathcal{F} \mid \diamond \mathcal{F} \mid \mathcal{F} \backslash \mathcal{F} \mid \mathcal{F} / \mathcal{F} \mid \mathcal{F} \bullet \mathcal{F}.$$

Figure 1 gives the (sequent-style) natural deduction presentation, together with the Curry-Howard term labelling¹. Judgements are of the form $\Gamma \vdash B$, with B a formula and Γ a structure term with formulas at the leaves. Antecedent structures are built according to the grammar $\mathcal{S} ::= \mathcal{F} \mid (\mathcal{S} \cdot \mathcal{S}) \mid \langle \mathcal{S} \rangle$. The binary structure-building operation $(- \cdot -)$ is the structural counterpart of the connective \bullet in the formula language. The unary structure-building operation $\langle - \rangle$ similarly is the counterpart of \diamond in the formula language.

With term labelling added, an antecedent term Γ with leaves $x_1 : A_1, \dots, x_n : A_n$ becomes a typing environment giving type declarations for the variables x_i . These variables constitute the parameters for the program t associated with the proof of the succedent type B . Intuitively, one can see a term-labeled proof as an algorithm to compute a meaning t of type B with parameters x_i of type A_i . In parsing a particular phrase, one substitutes the meaning of the constants (i.e. words) that make up for the parameters of this algorithm.

Notice that the term language respects the distinction between $/$ and \backslash : we use the ‘directional’ lambda terms of [19] with left versus right abstraction and application. The inference rules for \square and \diamond are reflected in the term language by \vee, \cup (Elimination) and \wedge, \cap (Introduction) respectively.

In addition to the logical rules for \diamond and \square , we are interested in formulating options for structural reasoning keyed to their presence. Consider the postulates expressed by the categorical morphisms of (1), or the corresponding inference rules of (2) in the N.D. format of Figure 1. These represent controlled forms of associativity and commutativity, explicitly licensed by the presence of \diamond (or its structural counterpart $\langle - \rangle$ in the sequent rules).

$$\diamond A \otimes (B \otimes C) \longrightarrow (\diamond A \otimes B) \otimes C \quad \diamond A \otimes (B \otimes C) \longrightarrow B \otimes (\diamond A \otimes C) \quad (1)$$

$$\frac{\Gamma[\langle (\Delta_1) \cdot \Delta_2 \rangle \cdot \Delta_3] \vdash t : B}{\Gamma[\langle \Delta_1 \rangle \cdot (\Delta_2 \cdot \Delta_3)] \vdash t : B} \text{Ass}_\diamond \quad \frac{\Gamma[\Delta_2 \cdot (\langle \Delta_1 \rangle \cdot \Delta_3)] \vdash t : B}{\Gamma[\langle \Delta_1 \rangle \cdot (\Delta_2 \cdot \Delta_3)] \vdash t : B} \text{Comm}_\diamond \quad (2)$$

Controlled forms of structural reasoning of this type have been used to model the dependencies between question words or relative pronouns and ‘gaps’ (physically unrealized hypothetical resources) that follow them. We illustrate with Dutch relative clauses, and refer the reader to [15] for a vector-based semantic analysis. Dutch, like Japanese, has verb-final word order in embedded clauses as show in (3a) which translates as (3b). Now consider the relative clause (3c). It has two possible interpretations, expressed by the translations (3d) and (3e). With a typing $(n \backslash n) / (np \backslash s)$ for the relative pronoun ‘die’ we can capture only the (3d) interpretation; the improved typing $(n \backslash n) / (\diamond \square np \backslash s)$ creates a derivational ambiguity that covers both the (3d) and the (3e) interpretation, where the

¹We restrict to the simply typed fragment, ignoring the \bullet operation.

$$\begin{array}{c}
\text{Terms: } \quad t, u ::= x \mid \lambda^r x.t \mid \lambda^l x.t \mid t \triangleleft u \mid u \triangleright t \mid t^{\cup} \mid t^{\cap} \mid t^{\vee} \mid t^{\wedge} \\
\\
\text{Typing rules:} \\
\\
\frac{}{x : A \vdash x : A} Ax \\
\\
\frac{\Gamma \cdot x : A \vdash t : B}{\Gamma \vdash \lambda^r x.t : B/A} I/ \quad \frac{x : A \cdot \Gamma \vdash t : B}{\Gamma \vdash \lambda^l x.t : A \setminus B} I \setminus \\
\\
\frac{\Gamma \vdash t : B/A \quad \Delta \vdash u : A}{\Gamma \cdot \Delta \vdash t \triangleleft u : B} E/ \quad \frac{\Gamma \vdash u : A \quad \Delta \vdash t : A \setminus B}{\Gamma \cdot \Delta \vdash u \triangleright t : B} E \setminus \\
\\
\frac{\langle \Gamma \rangle \vdash t : B}{\Gamma \vdash^{\wedge} t : \Box B} I \Box \quad \frac{\Gamma \vdash t : B}{\langle \Gamma \rangle \vdash^{\cap} t : \Diamond B} I \Diamond \\
\\
\frac{\Gamma \vdash t : \Box B}{\langle \Gamma \rangle \vdash^{\vee} t : B} E \Box \quad \frac{\Delta \vdash t : \Diamond A \quad \Gamma[(x:A)] \vdash u : B}{\Gamma[\Delta] \vdash u^{\cup} t/x : B} E \Diamond
\end{array}$$

Figure 1: $\mathbf{NL}\Diamond$. Proofs and terms. Antecedent structure terms must be non-empty. Notation $\Gamma[\Delta]$ for structure term Γ with substructure Δ .

latter relies on the ability of the $\Diamond\Box np$ hypothesis to ‘jump over’ the subject by means of $Comm_{\Diamond}$.

- a. (ik weet dat) $Bob_{np} Alice_{np}$ bewondert $_{np \setminus (np \setminus s)}$
- b. (I know that) Bob_{np} admires $_{(np \setminus s)/np}$ $Alice_{np}$
- c. man_n die $_{??}$ de_hond_{np} bijt $_{np \setminus (np \setminus s)}$ (3)
- d. man who bites the dog (= subject relativisation)
- e. man whom the dog bites (= object relativisation)

The crucial subderivations for the (3c) example schematically rely on the following steps (working upward): \setminus Introduction withdraws the $\Diamond\Box np$ hypothesis, \Diamond Elimination followed by zero or more steps of structural reasoning bring the hypothesis to the position where it can actually be used as a ‘regular’ np , thanks to the \Box Elimination proof of $\langle \Box np \rangle \vdash np$. The derived rule ($xleft$) in (4) telescopes this sequence of inference steps into a one-step inference, allowing for a succinct representation of the derivations.

$$\begin{array}{c}
\frac{z : \Box A \vdash z : \Box A}{\langle z : \Box A \rangle \vdash^{\vee} z : A} E \Box \\
\vdots \\
\frac{}{\Gamma[\langle z : \Box A \rangle \cdot \Delta] \vdash t : B} \\
\\
\frac{\vdots}{x : \Diamond \Box A \vdash x : \Diamond \Box A \quad \langle z : \Box A \rangle \cdot \Gamma[\Delta] \vdash t : B} (Ass_{\Diamond}, Comm_{\Diamond})^* \\
\frac{}{\Gamma[\Delta] \vdash \lambda^l x.t[\cup x/z] : \Diamond \Box A \setminus B} E \Diamond \\
\frac{x : \Diamond \Box A \cdot \Gamma[\Delta] \vdash t[\cup x/z] : B}{\Gamma[\Delta] \vdash \lambda^l x.t[\cup x/z] : \Diamond \Box A \setminus B} I \setminus \\
\frac{\Gamma[y : A \cdot \Delta] \vdash t : B}{\Gamma[\Delta] \vdash \lambda^l x.t[\vee \cup x/y] : \Diamond \Box A \setminus B} xleft
\end{array}
\tag{4}$$

Using our compiled inference rule, here are the derivations of both relativization readings, to be compared with those with the full uncompiled derivation in Appendix A. On the of the subject relativisation reading (3d), at the axioms, we show the constants (words) that will be substituted for the parameters of the proof term for the derivation. Also, in the structure terms on the left of the turnstile, we use these words instead of the parameter-type

pairs to enhance legibility. This derivation uses the $\diamond \square np$ hypothesis as the subject of the relative clause body; it simply relies on \diamond and \square Elimination, and doesn't involve structural reasoning.

$$\frac{\frac{\frac{\frac{\text{de}}{x_2 : np/n} \ell \quad \frac{\text{hond}}{y_2 : n} \ell}{\text{de} \cdot \text{hond} \vdash (x_2 \triangleleft y_2) : np} [E] \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{\frac{[_ \vdash x : np]^1 \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{(\text{de} \cdot \text{hond}) \cdot \text{bijt} \vdash ((x_2 \triangleleft y_2) \triangleright z_2) : np \setminus s} [\setminus E]} [\setminus E]} \quad \frac{\text{die}}{z_0 : (n \setminus n) / (\diamond \square np \setminus s)} \ell \quad \frac{[_ \vdash x : np]^1 \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{(\text{de} \cdot \text{hond}) \cdot \text{bijt} \vdash (x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) : s} [\setminus E]} [\setminus E]} \quad \frac{\frac{\text{man}}{y_0 : n} \ell \quad \frac{\text{die}}{z_0 : (n \setminus n) / (\diamond \square np \setminus s)} \ell \quad \frac{[_ \vdash x : np]^1 \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{(\text{de} \cdot \text{hond}) \cdot \text{bijt} \vdash \lambda^l x_1. (\vee^u x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) : \diamond \square np \setminus s} [xleft]^1}}{\text{die} \cdot ((\text{de} \cdot \text{hond}) \cdot \text{bijt}) \vdash (z_0 \triangleleft \lambda^l x_1. (\vee^u x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2))) : n \setminus n} [E]} [E]} \quad \frac{\text{man} \cdot (\text{die} \cdot ((\text{de} \cdot \text{hond}) \cdot \text{bijt})) \vdash (y_0 \triangleright (z_0 \triangleleft \lambda^l x_1. (\vee^u x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)))) : n} [E]} [E]$$

Contrast this with the derivation of the (3e) object relativisation interpretation. In this case the $\diamond \square np$ hypothesis is manoeuvred to the direct object position in the relative clause body thanks to the controlled commutativity option.

$$\frac{\frac{\frac{\frac{\text{de}}{x_2 : np/n} \ell \quad \frac{\text{hond}}{y_2 : n} \ell}{\text{de} \cdot \text{hond} \vdash (x_2 \triangleleft y_2) : np} [E] \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{\frac{[_ \vdash x : np]^1 \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{(\text{de} \cdot \text{hond}) \cdot \text{bijt} \vdash (x \triangleright z_2) : np \setminus s} [\setminus E]} [\setminus E]} \quad \frac{\text{die}}{z_0 : (n \setminus n) / (\diamond \square np \setminus s)} \ell \quad \frac{[_ \vdash x : np]^1 \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{(\text{de} \cdot \text{hond}) \cdot \text{bijt} \vdash (x_2 \triangleleft y_2) \triangleright (x \triangleright z_2) : s} [xleft]^1}}{\frac{\text{die} \cdot ((\text{de} \cdot \text{hond}) \cdot \text{bijt}) \vdash (z_0 \triangleleft \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee^u x_1 \triangleright z_2))) : \diamond \square np \setminus s} [E]} [E]} \quad \frac{\frac{\text{man}}{y_0 : n} \ell \quad \frac{\text{die}}{z_0 : (n \setminus n) / (\diamond \square np \setminus s)} \ell \quad \frac{[_ \vdash x : np]^1 \quad \frac{\text{bijt}}{z_2 : np \setminus (np \setminus s)} \ell}{(\text{de} \cdot \text{hond}) \cdot \text{bijt} \vdash \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee^u x_1 \triangleright z_2)) : \diamond \square np \setminus s} [xleft]^1}}{\text{die} \cdot ((\text{de} \cdot \text{hond}) \cdot \text{bijt}) \vdash (z_0 \triangleleft \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee^u x_1 \triangleright z_2))) : n \setminus n} [E]} [E]} \quad \frac{\text{man} \cdot (\text{die} \cdot ((\text{de} \cdot \text{hond}) \cdot \text{bijt})) \vdash (y_0 \triangleright (z_0 \triangleleft \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee^u x_1 \triangleright z_2)))) : n} [E]} [E]$$

Our aim in the following sections is to provide a compositional interpretation of the control operators and the structural reasoning licensed by them that allows us to simultaneously represent the co-existing interpretations of ambiguous utterances such as (3c).

3 Interpretation Spaces

Let us turn to the action of the interpretation homomorphism on the *types* of our extended Lambek calculus. In the approach introduced in [6], types are sent to density matrix spaces. These spaces are set up in a directionality-sensitive way, keeping in the semantics the distinction between left- or right-looking implications. Starting from the vector space V and its dual V^* , we use a modified Dirac notation to distinguish between two sets of basis of V , $\{|i'\rangle\}$ and $\{|^j\rangle\}$, and two sets of basis of V^* , $\{\langle^j|\}$ and $\{\langle i|\}$, obeying the orthogonality conditions

$$\langle i|i'\rangle = g_{ii'}, \quad \langle^j|^j\rangle = g^{jj}, \quad \langle i|^j\rangle = \delta_i^j, \quad \text{and} \quad \langle^j|\langle i| = \delta_i^j,$$

where g , called the *metric*, accounts for the eventual non-orthogonality of basis elements. In general, the basis vector $\langle^j|^i\rangle$ is obtained by conjugating and transposing the state $|^j\rangle$. When the basis is not orthogonal, this operation does not render the dual basis vector of $\langle^j|^i\rangle$, which by definition is orthogonal to it and in our notation is represented by $|i\rangle$, but instead gives another vector $\langle^j|^i\rangle$ that requires the metric tensor to describe its relation with it. Compare this with the case with only one set of basis for each space, obtained in the standard way: $\langle^j|^i\rangle$ coincides with $|i\rangle$ so that all basis vectors are orthogonal to each other, and there is therefore no need for a metric.

The *basic* building block for the interpretations is the density matrix space $\tilde{V} \equiv V \otimes V^*$. This space has density matrices as elements, which we will use as the starting representations of words, instead of vectors. Density matrices are 1) positive operators with 2) trace normalized to 1 [16]. In a physical system, this means that we can not only access the quantum properties of states, expressed as a linear combination of basis states of V or V^* , but we can also include the classical properties of a state, by constructing a basis of \tilde{V} and describing the states as any linear combination formed with these basis elements that obeys conditions 1) and 2). Because the range of representations is enlarged, their use has been proposed for linguistic applications [2,3,17,18], which we expand on here focusing on including the directionality of the calculus in this distributional representation. Defining the basis of V and V^* as we did before, we are able to construct a non-trivial basis for the density matrix space that carries over the structure of duality. For this space, we choose the basis formed by $|i\rangle$ tensored with $\langle i'|$, $\tilde{E} = \{|i\rangle\langle i'|\}$. We define the dual density matrix space $\tilde{V}^* \equiv V \otimes V^*$ and assign to the dual basis of this space the map that takes each basis element of \tilde{V} and returns a scalar. That basis is formed by $\langle j|$ tensored with $|j'\rangle$, $\tilde{D} = \{|j'\rangle\langle j|\}$, and is applied on the basis vectors of \tilde{V} via the trace operation

$$\text{Tr}(|i\rangle\langle i'|j'\rangle\langle j|) = \sum_l \langle l|i\rangle\langle i'|j'\rangle\langle j|l\rangle. \quad (5)$$

The *composite* spaces are formed via the binary operation \otimes (tensor product) and the unary operation $()^*$ (dual functor) that sends the elements of a density matrix basis to its dual basis, using the metric tensor. In the notation, we use \tilde{A} for density matrix spaces (basic or compound), and ρ , or subscripted $\rho_x, \rho_y, \rho_z, \dots \in \tilde{A}$ for elements of such spaces. The $()^*$ operation is involutive; it interacts with the tensor product as $(\tilde{A} \otimes \tilde{B})^* = \tilde{B}^* \otimes \tilde{A}^*$ and acts as identity on matrix multiplication.

The homomorphism that sends syntactic types to semantic spaces is the map $[\cdot]$. For primitive types it acts as

$$[s] = \tilde{S} \quad \text{and} \quad [np] = [n] = \tilde{N},$$

with S the vector space for sentence meanings, N the space for nominal expressions (common nouns, full noun phrases). For compound types we have

$$[A/B] = [A] \otimes [B]^* \quad \text{and} \quad [A \setminus B] = [A]^* \otimes [B].$$

This can be seen as an *operational* interpretation of formulae: a dualizing functor acting on one of the types, followed by a tensor product, also a functor, are identified with particular operations on elements, specifically by multiplying with the elements of a metric or by taking the trace².

We now turn to how to send the formulae decorated with unary modalities to semantic spaces, in a way that stays in this functorial/operational framework. Recall that in earlier work [14,15] modally marked formulae are interpreted in the same space as their undecorated versions, i.e. $[\diamond A] = [\square A] = [A]$.

To build a non-trivial interpretation of the unary connectives, we expand the interpretation space using the description of quantum states that distinguishes between their *spatial* and *spin* degrees of freedom. Let the $[\cdot]$ homomorphism give a description of the *spatial* components, encoding the numerically extracted distributional data. In addition to the spatial component, and commuting freely with the spatial parts, we introduce a new vector space, a density matrix space \mathfrak{S} with dimension $n \times n$ where the *spin* components are encoded. We denote this by the *n-level spin space*. Here we do not distinguish between covariant and contravariant components, making the standard Dirac notation the appropriate one to deal with this space. Accordingly, the basis is orthonormal and has elements in $\{|a\rangle\langle a'|\}$, with the values of a and a' ranging from 0 to $n-1$.

To obtain the full translation from syntactic types to their distributional interpretation spaces, we introduce an extended interpretation homomorphism that tensors the $[\cdot]$ interpretation of *all* types with a density matrix space \mathfrak{S} resulting in

$$[A] = [A] \otimes \mathfrak{S}. \quad (6)$$

²Equivalently, in a categorical distributional framework this corresponds to establishing a basis and taking either tensor contraction or multiplication as the operations that represent the η and ε maps at the element level.

For atoms and slash types, $\lceil \cdot \rceil$ stays as defined. For $\diamond A$ and $\square A$, we tensor $\lceil A \rceil$ with $\mathfrak{S} \otimes \mathfrak{S}^*$, the type for the matrix representation of the operators associated with \diamond and \square , that is,

$$\lceil \diamond A \rceil = \lceil \square A \rceil = \lceil A \rceil \otimes \mathfrak{S} \otimes \mathfrak{S}^*. \quad (7)$$

The key idea here is that by tensoring every type with an extra spin space via $\lceil \cdot \rceil$, the marked types have representations that encode maps from \mathfrak{S} to \mathfrak{S} coming from $\lceil \cdot \rceil$. This justifies the use of the same spin space to interpret the two markers, as they act as endomorphisms on the \mathfrak{S} space coming from $\lceil \cdot \rceil$, as in for lozenge $\lceil \diamond A \rceil = \lceil \diamond A \rceil \otimes \mathfrak{S}$ and similarly for box. At the *type* level, then, we find the structure to accommodate the operators $T_\diamond, T_\square \in \mathcal{L}(\mathfrak{S})$, for which the concrete distinct interpretations will then be provided at the *term* level. The key point of this structure is to give us precise control over the spin space as we interpret the unary modalities. Note that our connectives' interpretations do not interfere either with the distributional data that is stored in the spacial spaces, which is compatible with the interpretation of these connectives in previous work [14, 15]. The interpretations we assign to the unary connectives consist of operations that only modify elements of an ancillary space. By enlarging the distributional space with this new spin space, we can effectively find a distributional meaning for the unary connectives.

As an example, here is the $\lceil \cdot \rceil$ mapping for the relative pronoun type of (3c).

$$\begin{aligned} \lceil (n \setminus n) / (\diamond \square np \setminus s) \rceil &= \lceil (n \setminus n) / (\diamond \square np \setminus s) \rceil \otimes \mathfrak{S} \\ &= \lceil n \rceil^* \otimes \lceil n \rceil \otimes \lceil s \rceil^* \otimes \lceil np \rceil \otimes \underbrace{(\mathfrak{S} \otimes \mathfrak{S}^*)}_{T_\diamond} \otimes \underbrace{(\mathfrak{S} \otimes \mathfrak{S}^*)}_{T_\square} \otimes \mathfrak{S} \end{aligned} \quad (8)$$

4 Operational Interpretation of Lambek Rules

Given the new semantic spaces for the syntactic types, we now turn to the interpretation of the syntactic *derivations*, as encoded by their lambda proof terms, proving the soundness of the calculus presented in section 2 with respect to the semantics of section 3. In spin space, the operations that interpret different syntactic maps relate with the quantum postulates describing measurement and evolution of quantum systems [16].

Quantum measurement: Quantum measurements are described by a collection M_a of measurement operators, acting on the state space of the system being measured. The index a refers to the measurement outcomes that may occur in the experiment. If the state of the quantum system is ρ immediately before the measurement then the probability that result a occurs is given by $p(a) = \text{Tr}(M_a^\dagger M_a \rho)$ and the state of the system after the measurement is

$$\rho_a := \frac{M_a \rho M_a^\dagger}{p(a)}. \quad (9)$$

The measurement operators satisfy the completeness equation, $\sum_a M_a^\dagger M_a = I$. For an observable M with eigenvalues m and eigenvectors $|a\rangle$, a *projective* measurement is defined with $M_a = |a\rangle \langle a|$; in this context we say that a state has been projected onto $|a\rangle \langle a|$, and the quantum operator is then called a *projector*.

Evolution The evolution of a closed quantum system is described by a unitary transformation. That is, the state ρ^i of the system at time t_i is related to state ρ^{i+1} of the system at time t_{i+1} by a unitary operator U which depends only on these times. The state ρ^{i+1} relates with the previous one ρ^i by $\rho^{i+1} = U \rho^i U^\dagger$.

This correspondence is established via a function $\llbracket \cdot \rrbracket_g$ that associates each term t of type A with a semantic value, i.e. an element of $\lceil A \rceil$, the semantic space where meanings of type A live. For proof terms, $\llbracket \cdot \rrbracket$ is defined relative to an assignment function g ³, that provides a semantic value for the basic building blocks, viz. the variables that label

³The assignment function g should not be confused with the metric tensor g .

the axiom leaves of a proof, in this case independently for the spatial (S) and spin (\mathfrak{S}) components. A particular assignment $g_{x,kk'}^S$ is used to interpret the lambda abstraction in the spatial spaces:

Definition 4.1. The assignment $g_{x,kk'}^S$ is the assignment exactly like g^S except possibly for the parametric variable x of type A , which takes the value of the basis element of the interpreting space $|_k\rangle_{[A]}\langle_{k'}$.

The elements of the spin space are given by

$$\rho_x^{\mathfrak{S}} = \sum_{a,a'=0}^{n-1} \mathfrak{S} \mathbf{X}_{aa'} |a\rangle_{\mathfrak{S}} \langle a'|. \quad (10)$$

A pair of special assignment functions are $g_{x,I}^{\mathfrak{S}}$ and $g_{x,y}^{\mathfrak{S}}$, used to interpret the lambda abstraction in the spin space:

Definition 4.2. The assignment $g_{x,I}^{\mathfrak{S}}$ is the assignment exactly like $g^{\mathfrak{S}}$ except possibly for the parametric variable x of type A , which takes the value of the normalized identity, $I = \sum_a \frac{1}{\dim \mathfrak{S}} |a\rangle_{\mathfrak{S}} \langle a|$.

Definition 4.3. The assignment $g_{x,y}^{\mathfrak{S}}$ is the assignment exactly like $g^{\mathfrak{S}}$ except possibly for the parametric variable x of type A , which takes the value of variable y , also of type A .

The spatial interpretation of terms of types formed with binary connectives is as given in [6]. We reproduce here the main results, but focus on their interpretation in spin space. Further, we introduce the interpretation of the rules that introduce and eliminate unary connectives.

Some elimination rules will be interpreted in spin space using an instance of a projective measurement. Given a term u of type A and another term t of type B , we define a map $[[t^A]]_{g^{\mathfrak{S}}} * [[u^B]]_{g^{\mathfrak{S}}} : \mathfrak{S} \times \mathfrak{S} \rightarrow \mathfrak{S}$ acting on the interpretation of the terms in spin space:

$$[[t^A]]_{g^{\mathfrak{S}}} * [[u^B]]_{g^{\mathfrak{S}}} = \frac{\left([[u^B]]_{g^{\mathfrak{S}}} \right)^{\frac{1}{2}} \cdot [[t^A]]_{g^{\mathfrak{S}}} \cdot \left([[u^B]]_{g^{\mathfrak{S}}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{S}} \left(\left([[u^B]]_{g^{\mathfrak{S}}} \right)^{\frac{1}{2}} \cdot [[t^A]]_{g^{\mathfrak{S}}} \cdot \left([[u^B]]_{g^{\mathfrak{S}}} \right)^{\frac{1}{2}} \right)}, \quad (11)$$

with $(\cdot)^{\frac{1}{2}}$ such that when applied on an operator R we have that $(R)^{\frac{1}{2}} \cdot (R)^{\frac{1}{2}} = R$. Positive operators, such as density matrices, have a unique positive square root [1]. Physically, the spin split in its square-root acts as a measurement operator on the other input spin. Using normalization, the outcome is a well defined spin state ⁴.

4.1 Axiom

The axiom will be given by an element of the spatial spaces, tensored with an element of the spin space.

$$[[x^A]]_g = g(x^A) = \rho_x^{|A|} = [[x^A]]_{g^S} \otimes [[x^A]]_{g^{\mathfrak{S}}}, \quad (12)$$

where

$$[[x^A]]_{g^{\mathfrak{S}}} = \sum_{aa'} \mathfrak{S} \mathbf{X}_{aa'} |a\rangle_{\mathfrak{S}} \langle a'| \quad \text{and} \quad [[x^A]]_{g^S} = \sum_{ii'} {}^S \mathbf{X}^{ii'} |i\rangle_{[A]} \langle i'|. \quad (13)$$

⁴This is a generalization of one of the Frobenius algebras already used in [2] in the category $\mathbf{CPM}(\mathbf{FHilb})$, where, given the full density matrix representations of sentence, noun and verb, respectively $\rho(s)$, $\rho(n)$ and ρ , they relate by $\rho(s) = \rho(n)^{\frac{1}{2}} \rho(v) \rho(n)^{\frac{1}{2}}$. For a similar extension, see the definitions of "fuzz" and "phaser" in [3].

4.2 Introduction and elimination of binary connectives

Elimination of / and \

$$\llbracket (t \triangleleft u)^B \rrbracket_g \equiv \text{Tr}_{[A]} \left(\left(\llbracket t^{B/A} \rrbracket_{g^S} \cdot \llbracket u^A \rrbracket_{g^S} \right) \otimes \llbracket t^{B/A} \rrbracket_{g^\mathfrak{S}} * \llbracket u^A \rrbracket_{g^\mathfrak{S}} \right). \quad (14)$$

$$\llbracket (u \triangleright t)^B \rrbracket_g \equiv \text{Tr}_{[A]} \left(\left(\llbracket u^A \rrbracket_{g^S} \cdot \llbracket t^{A/B} \rrbracket_{g^S} \right) \otimes \llbracket u^A \rrbracket_{g^\mathfrak{S}} * \llbracket t^{A/B} \rrbracket_{g^\mathfrak{S}} \right). \quad (15)$$

Introduction of / and \

$$\llbracket (\lambda'_{x.t})^{B/A} \rrbracket_g \equiv \sum_{kk'} \left(\llbracket t^B \rrbracket_{g_{x,kk'}^S} \otimes |k'\rangle_{[A]^*} \langle k| \right) \otimes \llbracket t^B \rrbracket_{g_{x,l}^\mathfrak{S}}. \quad (16)$$

$$\llbracket (\lambda'_{x.t})^{A/B} \rrbracket_g \equiv \sum_{kk'} \left(|k'\rangle_{[A]^*} \langle k| \otimes \llbracket t^B \rrbracket_{g_{x,kk'}^S} \right) \otimes \llbracket t^B \rrbracket_{g_{x,l}^\mathfrak{S}}. \quad (17)$$

Syntactic equalities like beta reduction are interpreted as equalities in this model, as is shown in appendix D.

4.3 Introduction and elimination of unary connectives

The operation T_\diamond acting on elements of \mathfrak{S} is the linear combination of projectors T_\diamond^a onto pure states used as projectors $M_a = |a\rangle_{\mathfrak{S}} \langle a|$, generated by the eigenstates of an observable with n different eigenvalues, specified for a particular unary modality, indexed by $a \in \{0, \dots, n-1\}$. Applied on a state $\rho_x^\mathfrak{S}$, the general result is the mixed state

$$T_\diamond(\rho_x^\mathfrak{S}) = \sum_{a=0}^{n-1} c_a T_\diamond^a(\rho_x^\mathfrak{S}) \equiv \sum_{a=0}^{n-1} c_a (\rho_x^\mathfrak{S} * |a\rangle_{\mathfrak{S}} \langle a|) = \sum_{a=0}^{n-1} c_a \left(\frac{M_a \rho_x^\mathfrak{S} M_a}{\text{Tr}(M_a \rho_x^\mathfrak{S} M_a)} \right), \quad (18)$$

with $\sum_{a=0}^n c_a = 1$, $c_a \in \mathbb{R}$. Defining the ordering of the eigenstates by the increasing value of their corresponding index a , box elimination and diamond introduction will be interpreted as the projection onto the lowest eigenstate, effectively $c_0 = 1$ and $c_{a \neq 0} = 0$. By contrast, for the interpretation of axiom terms of types that include unary types, all coefficients can be non-zero.

The operation T_\square acts on elements by performing a unitary transformation, generated by the successive application of matrices $U_0 = \mathbb{1}$ and $U_b \in SU(n)$ on density matrices, for $b \in \{1, \dots, n^2 - 1\}$, represented as T_\square^b , for a particular representation and ordering. Again, for a state $\rho_x^\mathfrak{S}$, the application of this operation is

$$\left(T_\square^b(\rho_x^\mathfrak{S}) \right)^{d_b} = \begin{cases} \rho_x^\mathfrak{S} & \text{if } d_b = 0 \\ U_b \rho_x^\mathfrak{S} U_b^\dagger & \text{if } d_b = 1 \end{cases} \quad (19)$$

$$T_\square(\rho_x^\mathfrak{S}) = \left(T_\square^{n^2-1} \left(T_\square^{n^2-2} \left(\dots \left(T_\square^0(\rho_x^\mathfrak{S}) \right)^{d_0} \right) \right)^{d_{n^2-1}} \right)^{d_{n^2-1}} \quad (20)$$

where $()^\dagger$ indicates hermitian conjugation and $d_b \in \{0, 1\}$ ⁵. Diamond elimination and box introduction are interpreted as performing unitary transformation when $d_0 = 1$ and $d_{b \neq 0} \neq 0$.⁶

Elimination of \square : $\llbracket (\vee t)^B \rrbracket_g = \llbracket t^{\square B} \rrbracket_{g^S} \otimes T_\diamond^0 \left(\llbracket t^{\square B} \rrbracket_{g^\mathfrak{S}} \right)$

Elimination of \diamond :

$$\llbracket (u[\cup t/x])^B \rrbracket_g = \text{Tr}_{[A]} \left(\left[\llbracket t^{\diamond A} \rrbracket_{g^S} \cdot \sum_{kk'} \langle k' |_{[A]^*} \langle k | \otimes \llbracket u^B \rrbracket_{g_{x,kk'}^S} \right] \otimes T_\square^0 \left(\llbracket u^B \rrbracket_{g^\mathfrak{S}} \right) \right) \quad (21)$$

Introduction of \square and \diamond :

$$\llbracket (\wedge t)^{\square B} \rrbracket_g = \llbracket t^B \rrbracket_{g^S} \otimes T_\square^0 \left(\llbracket t^B \rrbracket_{g^\mathfrak{S}} \right), \quad \llbracket (\cap t)^{\diamond B} \rrbracket_g = \llbracket t^B \rrbracket_{g^S} \otimes T_\diamond^0 \left(\llbracket t^B \rrbracket_{g^\mathfrak{S}} \right) \quad (22)$$

4.4 Structural Reasoning

To interpret the derived inference rule, a *raising operator* S_+ performs a unitary transformation on the input state, which is a projected state since a diamond elimination rule has to be used to allow the application of controlled commutativity, and is applied as many times as nodes that need to be jumped to be in the right position to be extracted. We record that information by an index m on the substitution brackets of the proof term encoding the (*xleft*) inference. The index acts as a power on the raising operator, $(S_+)^m$, changing a state $\rho_a = |a\rangle_{\mathfrak{S}} \langle a|$ to $\rho_{a+m} = |a+m\rangle_{\mathfrak{S}} \langle a+m|$, where we use the convention that a matrix to the zeroth power is the identity matrix. Additionally the derived inference rule is interpreted using the previously given interpretations of \square and \diamond .

Derived Inference Rule

xleft: Premise t^B with subterm y^A at location m ; conclusion $(\lambda^l x.t[\vee \cup x/y]^m)^{\diamond \square A \setminus B}$:

$$\begin{aligned} & \llbracket (\lambda^l x.t[\vee \cup x/y]^m)^{\diamond \square A \setminus B} \rrbracket_g = \\ & = \sum_{l'l'} \langle l' |_{[A]^*} \langle l | \otimes \left[\text{Tr}_{[A]} \left(\left[\llbracket x^{\diamond \square A} \rrbracket_{g^S} \cdot \sum_{kk'} \langle k' |_{[A]^*} \langle k | \otimes \llbracket t^B \rrbracket_{g_{y,kk'}^S} \right] \right) \right]_{g_{x,l'l'}^S} \end{aligned} \quad (23)$$

$$\otimes \left[\llbracket t^B \rrbracket_{g_{y,l}^\mathfrak{S}} * T_\square^0 \left((S_+)^m T_\diamond^0 \left(\llbracket x^{\diamond \square A} \rrbracket_{g^\mathfrak{S}} \right) \left((S_+)^{\dagger} \right)^m \right) \right]_{g_{x,l}^\mathfrak{S}} \quad (24)$$

5 Two-level spin space

The structural ambiguity at hand will be treated using a two-level spin space, since we have two ambiguous readings. This space is used to encode spin states of fermionic particles, with spin 1/2, such as electrons and protons. A helpful geometric visualization of the states in this space is the *Bloch sphere*, in fig. 2.

⁵Equation 20 can possibly be extended with permutations over the order of application of T_\square^b .

⁶Looking at the adjoint properties of the unary connectives, the implications $\diamond \square A \rightarrow A \rightarrow \square \diamond A$ are consistent with the non-commutativity of the operations that interpret these connectives, measurement and evolution, and also with the relationships between the sets obtained via these operations.

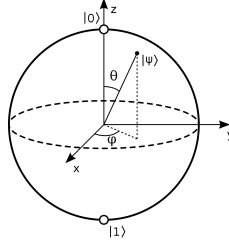


Figure 2: Bloch sphere representation of a two-level quantum state, also called a *qubit*. The general form of a state on the surface is $|\Psi\rangle = \left(\cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle\right) e^{i\gamma}$. The global phase $e^{i\gamma}$ is not represented because it has no effect on the density matrix. A product of states $\rho^{\text{pure}} = |\Psi\rangle\langle\Psi|$ is called a *pure state*, represented on the surface of the sphere. Otherwise the states are called *mixed states* and live inside of the sphere.

To interpret the action of the unary connectives in the spin space, we suppose that the particles with spin, our words in this case, are subjected to a uniform magnetic field pointing in the z direction. Using natural units, let $S_z = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ be the spin operator in the z direction. The eigenvectors of this operator are the orthogonal states $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, using the standard matrix representation. On the Bloch sphere, these states correspond to the north and south poles, respectively. The corresponding eigenvalues are $e_0 = -\frac{1}{2}$ and $e_1 = \frac{1}{2}$. This is the operator that we will use to interpret our unary modality. Thus T_\diamond is the set formed by linear combinations of states $\rho_0 = |0\rangle\langle 0|$ and $\rho_1 = |1\rangle\langle 1|$, the states that lie on the z -axis inside the Bloch sphere.

The raising operator is $S_+ = S_x + iS_y = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$. Once applied on ρ_0 the result is ρ_1 , and a further application has a null result.

6 Going Dutch again

To illustrate the interpretation process, we return to our Dutch relative clause example "man die de hond bijt", and show how we handle the derivational ambiguity. The lexicon below has the syntactic type assignments and the corresponding semantic spaces:

	syn type A	$[A]$
man'	n	$\tilde{N} \otimes \mathfrak{S}$,
die'	$(n \setminus n) / (\diamond \square np \setminus s)$	$\tilde{N}^* \otimes \tilde{N} \otimes \tilde{S}^* \otimes \tilde{N} \otimes (\mathfrak{S} \otimes \mathfrak{S}^*) \otimes (\mathfrak{S} \otimes \mathfrak{S}^*) \otimes \mathfrak{S}$,
de hond'	np	$\tilde{N} \otimes \mathfrak{S}$,
bijt'	$np \setminus np \setminus s$	$\tilde{N}^* \otimes \tilde{N}^* \otimes \tilde{S} \otimes \mathfrak{S}$.

In order to compute the interpretations given by the two above derivations, we start from the following primitive interpretations:

$$\llbracket \text{man}^{,n} \rrbracket_I = \sum_{rr',ii'} S \mathbf{M}^{rr'} |r\rangle_{[N]} \langle r'| \otimes \mathfrak{S} \mathbf{M}_{ii'} |i\rangle \langle i'|, \quad (25)$$

$$\llbracket \text{die}^{,(n \setminus n) / (\diamond \square np \setminus s)} \rrbracket_I = \sum_{kk',ll',mm',nn',ii'} S \mathbf{D}_{kk'm'm'}^{ll'nn'} |k' m' \rangle_{[N]^* \otimes [N] \otimes ([S]^* \otimes [N])} \langle l' n' | \otimes \mathfrak{S} \mathbf{D}_{ii'} |i\rangle \langle i'|; \quad (26)$$

$$\llbracket \text{de hond}^{,np} \rrbracket_I = \sum_{jj',ii'} S \mathbf{H}^{jj'} |j\rangle_{\tilde{N}} \langle j'| \otimes \mathfrak{S} \mathbf{H}_{ii'} |i\rangle \langle i'|; \quad (27)$$

$$\llbracket \text{bijt}^{,np \setminus np \setminus s} \rrbracket_I = \sum_{oo',pp',qq',ii'} S \mathbf{B}_{o'o,p'p}^{qq'} |o' p' \rangle_{[N]^* \otimes [N]^* \otimes [S]} \langle q' | \otimes \mathfrak{S} \mathbf{B}_{ii'} |i\rangle \langle i'|. \quad (28)$$

To obtain the correct contractions in the spatial components, that are related either to the subject or object relativization readings, the role of the hypothesis x is crucial: interpreted as in eq.13, it contracts with the interpretation of "bijt" as the interpretations of the slash elimination rules prescribe, either in subject or object position. Its most important role is in the latter, blocking "de hond" from taking the immediate object position contraction. After that, variable x is extracted using the *xleft* rule, in a way that keeps all the other contractions unchanged, and keeping the right form such that "die" can contract in the correct position.

With respect to the spin components, note that the spin values in eqs.25, 27 and 28 could be any point inside the Bloch sphere, but by eqs.8, 18 and 19, the spin value in eq.26 must live on the z -axis, that is, be represented by a diagonal matrix. For the instantiations of all of these words, the only requirement is that none has the value of ρ_0 or ρ_1 . Then, following the application of the rules using these values, the derived inference rule will produce a spin state equal to ρ_0 or ρ_1 , which will be preserved as it interacts with other spin states via slash elimination. The detailed calculations that interpret the different readings are in Appendix B and schematic representations of the contractions can be found in figs.3 and 4.

For the first reading the result is

$$\llbracket \text{man_die_de_hond_bijt} \rrbracket_I^1 = \sum_{rr',ll',mm',nn'} S \mathbf{M}^{rr'} S \mathbf{D}_{r'r'm'm'}^{ll'nn'} S \mathbf{H}^{jj'} S \mathbf{B}_{j',n'n}^{mm'} |l\rangle_{[N]} \langle l' | \otimes |0\rangle \langle 0|, \quad (29)$$

and for the second reading it is

$$\llbracket \text{man_die_de_hond_bijt} \rrbracket_I^2 = \sum_{rr',ll',mm',nn'} S \mathbf{M}^{rr'} S \mathbf{D}_{r'r'm'm'}^{ll'nn'} S \mathbf{H}^{jj'} S \mathbf{B}_{n'n,j'j}^{mm'} |l\rangle_{[N]} \langle l' | \otimes |1\rangle \langle 1|. \quad (30)$$

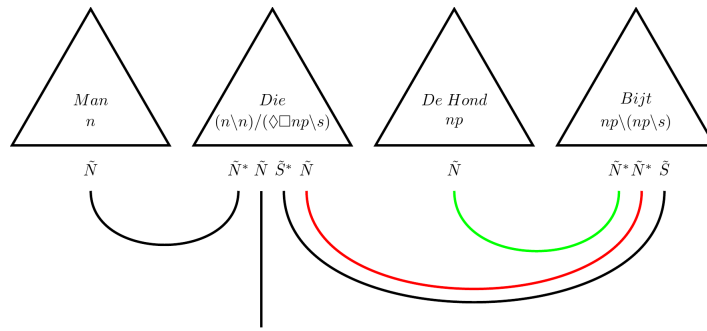


Figure 3: Representation of spatial contractions corresponding to the subject relativisation reading of "man die de hond bijt", according to eq.29, tensored with the north-pole state on the Bloch sphere.

The final interpretation of the ambiguous phrase is given by the direct sum of the two unambiguous interpretations, weighted by parameters p_1 and p_2 that express the likelihood of each reading:

$$\llbracket \text{man_die_de_hond_bijt} \rrbracket_I = p_1 \llbracket \text{man_die_de_hond_bijt} \rrbracket_I^1 \oplus p_2 \llbracket \text{man_die_de_hond_bijt} \rrbracket_I^2. \quad (31)$$

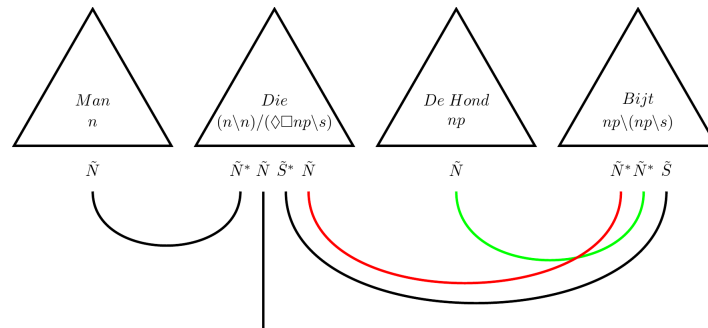


Figure 4: Representation of contractions corresponding to the object relativisation reading of "man die de hond bijt", according to eq.30, tensored with the south-pole state on the Bloch sphere.

7 Discussion and Conclusion

In this extended abstract we extended the interpretation space with a spin degree of freedom, showing how that can preserve extra information about the proof. We showed how interpreting the meanings of words directly as density matrices introduces a framework that can be used to encode higher-level content. Open questions that we want to address in future work have to do with how these changes scale up when we have more than one unary modality, possibly by associating them with eigenstates of different operators, or when more than two ambiguous readings are possible, using a larger spin space and an appropriate raising operator. Interesting too is to relate our approach, where lambda terms are directly interpreted using elements and operations over them, with Kripke frames on vector spaces [7], defining the valuation sets with the accessibility relations that translate into our operations, unveiling a stronger connection with the logic of residuation. Also relevant would be to compare our take on interpreting certain logic connectives using quantum mechanical operations with the mirror field of quantum logic [4] that aims at interpreting quantum mechanics using logic tool, particularly modal logic [4] which is at the root of our unary connectives, where too an association between projections and the logic of possibility (\diamond in our notation) is suggested. Finally, further research will have to show how the probability coefficients can be extracted from derivational data, and whether it is possible to go from the subject relativization reading to the object relativization reading applying only permutation operators as is done in [6] for syntactic ambiguities and, in that case, what the connection with the derivation precisely is.

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A.3 Formal semantics of relative pronouns

To obtain the usual ‘formal semantics’ terms, one substitutes for the parameter z_0 the lexical program for the word ‘die’:

$$\text{DIE} = \lambda x \lambda y \lambda z. ((y z) \wedge (x \overset{\cap}{\wedge} z))$$

which then, after β conversion and cap-cup and wedge-vee cancellation, reduces to

$$\lambda z. ((\text{MAN } z) \wedge ((\text{BIJT (DE HOND)) } z)) \quad (\text{subject reading})$$

$$\lambda z. ((\text{MAN } z) \wedge ((\text{BIJT } z) (\text{DE HOND}))) \quad (\text{object reading})$$

B Concrete interpretation of relative clauses

The derivations in 2 have a final term that depends on the variables y_0 , z_0 , x_2 , y_2 , z_2 and x_1 . The latter is a *bound* variable (as well as the intermediate variable x), due to the lambda abstraction, and the former are *free* variables. Bound variables can be substituted by any free variable during the derivation, via beta reduction, and will take the value of that variable, contrasting with free variables that will be substituted by constants, and interpreted accordingly. An assignment function g assigns bound variables to a later-to-be-defined constant, and assigns free variables to specific constants, here our words. In our assignment, taken as an example, the assignment function gives $g(y_0) = \text{man}'$ but $g(x_1)$ remain in this form, until x_1 is substituted by a free variable. Alternatively we can represent the free variables as bound variables using a lambda abstraction, applied on a constant: $\lambda y_0. y_0(\text{man}') \rightarrow \text{man}'$.

Looking at the interpretation of any variable stated in the interpretation of the axiom rule in eq.13 and comparing with the interpretation of the constants in eqs.25 to 28, we note that both represent the density matrix entries in a symbolic form, where we can apply directly operations like trace and matrix multiplication in the spatial components, or spin operators in the spin components. This permits that, when we perform these calculation step by step using each rule, we can perform them directly on the symbolic representations of interpretations of constants, in eqs.25 to 28, as well as of variables that naturally take the same form as states in eq.13, since it can potentially take the value of *any* other constant.

Therefore, one can impose an assignment that will interpret our particular Dutch relative clause "man die de hond bijt" g that instantiates the free variables like so:

$$\llbracket (x_2 \triangleleft y_2) \rrbracket_g = \llbracket \text{de_hond}^{np} \rrbracket_I, \quad (32)$$

$$\llbracket (z_2) \rrbracket_g = \llbracket \text{bijt}^{np \setminus np \setminus s} \rrbracket_I, \quad (33)$$

$$\llbracket (z_0) \rrbracket_g = \llbracket \text{die}^{(n \setminus) / (np \setminus s)} \rrbracket_I, \quad (34)$$

$$\llbracket (y_0) \rrbracket_g = \llbracket \text{man}^{n'} \rrbracket_I \quad (35)$$

and instantiates the bound variable x according to eq.13.

Substituting these directly in the derivations, we can, step by step, arrive at the final different readings. In what follows we give a full breakdown of these steps, splitting between spatial and spin components, and between subject and object relativization.

B.1 Interpretations in $\llbracket \cdot \rrbracket$:

B.1.1 Subject Relativization

The interpretation of this derivation starts by making use of the interpretation of E_{\setminus} as given in eq.15, substituting the variables by the assigned constants as described above.

$$\begin{aligned}
\llbracket (x_2 \triangleleft y_2) \triangleright z_2 \rrbracket_{g^S} &= \text{Tr}_{\tilde{N}} \left(\llbracket (x_2 \triangleleft y_2) \rrbracket_{g^S} \cdot \llbracket z_2 \rrbracket_{g^S} \right) \\
&= \text{Tr}_{\tilde{N}} \left(\sum_{jj'} {}^S \mathbf{H}^{jj'} |j\rangle_{\tilde{N}} \langle j'| \cdot \sum_{oo', pp', qq'} {}^S \mathbf{B}_{o'o, p'p}^{qq'} |o'p'q\rangle_{[N]^* \otimes [N]^* \otimes [S]} \langle op'q| \right) \\
&= \sum_{jj', pp', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{j'j, p'p}^{qq'} |p'q\rangle_{[N]^* \otimes [S]} \langle p'q|
\end{aligned} \tag{36}$$

Then we use again eq.15 and interpret the variable x using axiom rule as in eq.13.

$$\begin{aligned}
\llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g^S} &= \text{Tr}_{\tilde{N}} \left(\llbracket x \rrbracket_{g^S} \cdot \llbracket (x_2 \triangleleft y_2) \triangleright z_2 \rrbracket_{g^S} \right) \\
&= \text{Tr}_{\tilde{N}} \left(\sum_{ii'} {}^S \mathbf{X}^{ii'} |i\rangle_{\tilde{N}} \langle i'| \cdot \sum_{jj', pp', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{j'j, p'p}^{qq'} |p'q\rangle_{[N]^* \otimes [S]} \langle p'q| \right) \\
&= \sum_{ii', jj', qq'} {}^S \mathbf{X}^{ii'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{j'j, i'i}^{qq'} |q\rangle_{[S]} \langle q'|
\end{aligned} \tag{37}$$

To use the *xleft* rule, we first interpret the previous term in the assignment $g_{x, ll'}^S$, as described in Def.4.1. recalculating the previous interpretation using the basis of its interpretation space instead of eq.13.

$$\begin{aligned}
\llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g_{x, ll'}^S} &= \text{Tr}_{\tilde{N}} \left(|l\rangle_{[N]} \langle l'| \cdot \llbracket (x_2 \triangleleft y_2) \triangleright z_2 \rrbracket_{g^S} \right) \\
&= \sum_{jj', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{j'j, i'l}^{qq'} |q\rangle_{[S]} \langle q'|
\end{aligned} \tag{38}$$

We simplify the spatial interpretation of *xleft* as given in eq.23, using that x and y are interpreted both interpreted in $[A]$, since $\llbracket \diamond \square A \rrbracket = [A]$:

$$\begin{aligned}
&\llbracket (\lambda^l x. t[\vee \cup x/y]^m) \diamond \square A \setminus B \rrbracket_{g^S} = \\
&= \sum_{ll'} |l'\rangle_{[A]^*} \langle l| \otimes \left[\text{Tr}_{[A]} \left(\llbracket x \diamond \square A \rrbracket_{g^S} \cdot \sum_{kk'} |k'\rangle_{[A]^*} \langle k| \otimes \llbracket t^B \rrbracket_{g_{y, kk'}^S} \right) \right]_{g_{x, ll'}^S} \\
&= \sum_{ll'} |l'\rangle_{[A]^*} \langle l| \otimes \text{Tr}_{[A]} \left(|l\rangle_{[A]} \langle l'| \cdot \sum_{kk'} |k'\rangle_{[A]^*} \langle k| \otimes \llbracket t^B \rrbracket_{g_{y, kk'}^S} \right) \\
&= \sum_{ll'} |l'\rangle_{[A]^*} \langle l| \otimes \llbracket t^B \rrbracket_{g_{y, ll'}^S}.
\end{aligned} \tag{39}$$

Using this simplified form, we see that multiplying with the dual basis of the space that interprets both x and x_1 results in an expression that will take any value of a variable of that type, precisely the goal of the lambda abstraction.

$$\begin{aligned}
\llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g^s} &= \sum_{l'} |l'\rangle_{[N]^*} \langle l'| \otimes \llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g^s, l'}^s \\
&= \sum_{l'} |l'\rangle_{[N]^*} \langle l'| \otimes \sum_{jj', qq'} S \mathbf{H}^{jj'} S \mathbf{B}_{j'j, l'l'}^{qq'} |q\rangle_{[S]} \langle q'| \\
&= \sum_{l', jj', qq'} S \mathbf{H}^{jj'} S \mathbf{B}_{j'j, l'l'}^{qq'} |l'\rangle_{[N]^* \otimes [S]} \langle l' q'|
\end{aligned} \tag{40}$$

To finalize, the next two steps consist in the application of the interpretations of $E_/\mathit{}$ (eq.14) and E_\backslash (eq.15), respectively, resulting in the spatial part of eq. 29.

$$\begin{aligned}
\llbracket z_0 \triangleleft \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g^s} &= \text{Tr}_{\bar{S}} \left(\text{Tr}_{\bar{N}} \left(\llbracket z_0 \rrbracket_{g^s} \cdot \llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g^s} \right) \right) \\
&= \text{Tr}_{\bar{S}} \left(\text{Tr}_{\bar{N}} \left(\sum_{kk', tt', mm', nn'} S \mathbf{D}_{k'l'k'm'l'm}^{tt' nm'} |k' m'\rangle_{[N]^* \otimes [N] \otimes ([S]^* \otimes [N])} \langle k' m' t' n' | \cdot \sum_{l', jj', qq'} S \mathbf{H}^{jj'} S \mathbf{B}_{j'j, l'l'}^{qq'} |l' q\rangle_{[N]^* \otimes [S]} \langle l' q'| \right) \right) \\
&= \sum_{kk', tt', mm', nn', jj'} S \mathbf{D}_{k'l'k'm'l'm}^{tt' nm'} S \mathbf{H}^{jj'} S \mathbf{B}_{j'j, n'n}^{mm'} |k' t'\rangle_{[N]^* \otimes [N]} \langle k' t' |
\end{aligned} \tag{41}$$

$$\begin{aligned}
\llbracket y_0 \triangleright (z_0 \triangleleft \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g^s} &= \text{Tr}_{\bar{N}} \left(\llbracket y_0 \rrbracket_{g^s} \cdot \llbracket z_0 \triangleleft \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g^s} \right) \\
\text{Tr}_{\bar{N}} \left(\sum_{rr'} S \mathbf{M}^{rr'} |r\rangle_{[N]} \langle r'| \cdot \sum_{kk', tt', mm', nn', jj'} S \mathbf{D}_{k'l'k'm'l'm}^{tt' nm'} S \mathbf{H}^{jj'} S \mathbf{B}_{j'j, n'n}^{mm'} |k' t'\rangle_{[N]^* \otimes [N]} \langle k' t' | \right) \\
&= \sum_{rr', tt', mm', nn', jj'} S \mathbf{M}^{rr'} S \mathbf{D}_{r'r, m'l'm}^{tt' nm'} S \mathbf{H}^{jj'} S \mathbf{B}_{j'j, n'n}^{mm'} |t\rangle_{[N]} \langle r'| \\
&= \llbracket \text{man_die_de_hond_bijt} \rrbracket_{\bar{S}}^1
\end{aligned} \tag{42}$$

$$= \llbracket \text{man_die_de_hond_bijt} \rrbracket_{\bar{S}}^1 \tag{43}$$

B.1.2 Object relativization

This derivation is very similar to the previous, except that on the first application of E_\backslash the bound variable x is introduced as the argument of z_2 , and only on the next application of the rule is $(x_2 \triangleleft y_2)$ taken as an argument.

$$\begin{aligned}
\llbracket x \triangleright z_2 \rrbracket_{g^s} &= \text{Tr}_{\bar{N}} \left(\llbracket x \rrbracket_{g^s} \cdot \llbracket z_2 \rrbracket_{g^s} \right) \\
&= \text{Tr}_{\bar{N}} \left(\sum_{ii'} S \mathbf{X}^{ii'} |i\rangle_{\bar{N}} \langle i'| \cdot \sum_{oo', pp', qq'} S \mathbf{B}_{o'o, p'p}^{qq'} |o' p' q\rangle_{[N]^* \otimes [N]^* \otimes [S]} \langle o' p' q' | \right) \\
&= \sum_{ii', pp', qq'} S \mathbf{X}^{ii'} S \mathbf{B}_{i'i, p'p}^{qq'} |p' q\rangle_{[N]^* \otimes [S]} \langle p' q' |
\end{aligned} \tag{44}$$

$$\begin{aligned}
\llbracket (x_2 \triangleleft y_2) \triangleright (x \triangleright z_2) \rrbracket_{g^s} &= \text{Tr}_{\bar{N}} \left(\llbracket (x_2 \triangleleft y_2) \rrbracket_{g^s} \cdot \llbracket x \triangleright z_2 \rrbracket_{g^s} \right) = \text{Tr}_{\bar{N}} \left(\llbracket (x_2 \triangleleft y_2) \rrbracket_{g^s} \cdot \text{Tr}_{\bar{N}} \left(\llbracket x \rrbracket_{g^s} \cdot \llbracket z_2 \rrbracket_{g^s} \right) \right) \\
&= \text{Tr}_{\bar{N}} \left(\sum_{jj'} S \mathbf{H}^{jj'} |j\rangle_{\bar{N}} \langle j'| \cdot \sum_{ii', pp', qq'} S \mathbf{X}^{ii'} S \mathbf{B}_{i'i, p'p}^{qq'} |p' q\rangle_{[N]^* \otimes [S]} \langle p' q' | \right) \\
&= \sum_{jj', ii', qq'} S \mathbf{H}^{jj'} S \mathbf{X}^{ii'} S \mathbf{B}_{i'i, j'j}^{qq'} |q\rangle_{[S]} \langle q'|
\end{aligned} \tag{45}$$

Note at this point that, due to changing the ordering of contraction, when compared with the subject relativization reading, the matrix indices are contracted differently from eq.37. We see now what the role of the hypotheses x is: to block $(x_2 \triangleright y_2)$ from contracting inevitably as the first argument of z_2 . Now that the contraction is in line with what we want for an object relativization reading, we will extract variable x via *xleft*. To do that, we first reinterpret the previous term using the assignment $g_{x, ll'}^S$. To substitute the interpretation of x by that of its basis elements we need to go further into de proof, when compared with the subject relativization reading.

$$\begin{aligned}
\llbracket (x_2 \triangleleft y_2) \triangleright (x \triangleright z_2) \rrbracket_{g_{x, ll'}^S} &= \text{Tr}_{\bar{N}} \left(\llbracket (x_2 \triangleleft y_2) \rrbracket_{g^S} \cdot \text{Tr}_{\bar{N}} \left(|l\rangle_{[N]} \langle l'| \cdot \llbracket z_2 \rrbracket_{g^S} \right) \right) \\
&= \text{Tr}_{\bar{N}} \left(\sum_{jj'} {}^S \mathbf{H}^{jj'} |j\rangle_{\bar{N}} \langle j'| \cdot \sum_{pp', qq'} {}^S \mathbf{B}_{l'l, p'p}^{qq'} |p'\rangle_{[N]^* \otimes [S]} \langle q'| \right) \\
&= \sum_{jj', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{l'l, j'j}^{qq'} |q\rangle_{[S]} \langle q'|
\end{aligned} \tag{46}$$

The following steps are as before, with the final result referring to eq.30.

$$\begin{aligned}
\llbracket \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee x_1 \triangleright z_2)) \rrbracket_{g^S} &= \sum_{ll'} |l'\rangle_{[N]^*} \langle l| \otimes \llbracket ((x_2 \triangleleft y_2) \triangleright (x \triangleright z_2)) \rrbracket_{g_{x, ll'}^S} \\
&= \sum_{ll'} |l'\rangle_{[N]^*} \langle l| \otimes \sum_{jj', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{l'l, j'j}^{qq'} |q\rangle_{[S]} \langle q'|
\end{aligned} \tag{47}$$

$$= \sum_{ll', jj', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{l'l, j'j}^{qq'} |l'\rangle_{[N]^* \otimes [S]} \langle l| \tag{48}$$

$$\begin{aligned}
\llbracket z_0 \triangleleft \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee x_1 \triangleright z_2)) \rrbracket_{g^S} &= \text{Tr}_{\bar{S}} \left(\text{Tr}_{\bar{N}} \left(\llbracket z_0 \rrbracket_{g^S} \cdot \llbracket \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee x_1 \triangleright z_2)) \rrbracket_{g^S} \right) \right) \\
&= \text{Tr}_{\bar{S}} \left(\text{Tr}_{\bar{N}} \left(\sum_{kk', tt', mm', nn'} {}^S \mathbf{D}_{k'k m'm}^{tt' nn'} |k'\rangle_{[N]^* \otimes [N] \otimes ([S]^* \otimes [N])} \langle k m n' | \cdot \sum_{ll', jj', qq'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{l'l, j'j}^{qq'} |l'\rangle_{[N]^* \otimes [S]} \langle l q' | \right) \right) \\
&= \sum_{kk', tt', mm', nn', jj'} {}^S \mathbf{D}_{k'k m'm}^{tt' nn'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{n'n, j'j}^{mm'} |k'\rangle_{[N]^* \otimes [N]} \langle k t' |
\end{aligned} \tag{49}$$

$$\begin{aligned}
\llbracket y_0 \triangleright (z_0 \triangleleft \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee x_1 \triangleright z_2))) \rrbracket_{g^S} &= \text{Tr}_{\bar{N}} \left(\llbracket y_0 \rrbracket_{g^S} \cdot \llbracket z_0 \triangleleft \lambda^l x_1. ((x_2 \triangleleft y_2) \triangleright (\vee x_1 \triangleright z_2)) \rrbracket_{g^S} \right) \\
\text{Tr}_{\bar{N}} \left(\sum_{rr'} {}^S \mathbf{M}^{rr'} |r\rangle_{[N]} \langle r'| \cdot \sum_{kk', tt', mm', nn', jj'} {}^S \mathbf{D}_{k'k m'm}^{tt' nn'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{n'n, j'j}^{mm'} |k'\rangle_{[N]^* \otimes [N]} \langle k t' | \right) \\
&= \sum_{rr', tt', mm', nn', jj'} {}^S \mathbf{M}^{rr'} {}^S \mathbf{D}_{r'r m'm}^{tt' nn'} {}^S \mathbf{H}^{jj'} {}^S \mathbf{B}_{n'n, j'j}^{mm'} |t\rangle_{[N]} \langle r'| \\
&= \llbracket \text{man_die_de_hond_bijt} \rrbracket_{\bar{S}}^2.
\end{aligned} \tag{50}$$

B.2 Interpretations in \mathfrak{S} :

B.2.1 Subject Relativization

We start by using the interpretations of variables in the interpretation of E_{\setminus} as given in eq. 15, which are particular forms of eq. 11. The variables can have any value with the only requirement that it is neither ρ_0 nor ρ_1 . This is

because the resulting states must have a non-zero probability of being projected on either of these states, which is necessary for the following step.

$$\llbracket (x_2 \triangleleft y_2) \triangleright z_2 \rrbracket_{g\mathfrak{E}} = \llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} = \frac{\left(\llbracket z_2 \rrbracket_{g\mathfrak{E}}\right)^{\frac{1}{2}} \cdot \llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket z_2 \rrbracket_{g\mathfrak{E}}\right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left(\left(\llbracket z_2 \rrbracket_{g\mathfrak{E}}\right)^{\frac{1}{2}} \cdot \llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket z_2 \rrbracket_{g\mathfrak{E}}\right)^{\frac{1}{2}} \right)}. \quad (52)$$

$$\begin{aligned} \llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g\mathfrak{E}} &= \llbracket x \rrbracket_{g\mathfrak{E}} * \left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right) \\ &= \frac{\left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \left(\llbracket x \rrbracket_{g\mathfrak{E}} \right) \cdot \left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left(\left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \left(\llbracket x \rrbracket_{g\mathfrak{E}} \right) \cdot \left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \right)} \end{aligned} \quad (53)$$

Looking at the interpretation of $xleft$ in eq. 23, first the previous state is projected onto ρ_0 using T_{\diamond}^0 . Then, since controlled commutativity is not used, $m = 0$ and $(S_+)^0 = \mathbb{1}$. Lastly, we recall that in our definitions $U_0 = \mathbb{1}$.

$$T_{\diamond}^0 \left(\llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g\mathfrak{E}} \right) = \llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g\mathfrak{E}} * |0\rangle_{\mathfrak{E}} \langle 0| = |0\rangle_{\mathfrak{E}} \langle 0| \quad (54)$$

$$\begin{aligned} \llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} &= \\ &= T_{\square}^0 \left((S_+)^0 T_{\diamond}^0 \left(\llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g\mathfrak{E}} \right) \left((S_+)^{\dagger} \right)^0 \right) \\ &= |0\rangle_{\mathfrak{E}} \langle 0| \end{aligned} \quad (55)$$

In the following two steps, the interpretations of rules $E_{/}$ and E_{\setminus} are used. Pure state ρ_0 will be preserved, taking into account that

$$\llbracket t^A \rrbracket_{g\mathfrak{E}} * \llbracket u^B \rrbracket_{g\mathfrak{E}} = \llbracket u^B \rrbracket_{g\mathfrak{E}}, \quad (56)$$

when $\llbracket u^B \rrbracket_{g\mathfrak{E}}$ equals ρ_0 or ρ_1 . To show this, take $\llbracket t^A \rrbracket_{g\mathfrak{E}} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ and $\llbracket u^B \rrbracket_{g\mathfrak{E}} = |0\rangle_{\mathfrak{E}} \langle 0| = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$,

$$\llbracket t^A \rrbracket_{g\mathfrak{E}} * \llbracket u^B \rrbracket_{g\mathfrak{E}} = \frac{\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}}{\text{Tr} \left(\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \right)} = \frac{\begin{pmatrix} 0 & 0 \\ 0 & d \end{pmatrix}}{d} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}, \quad (57)$$

and similarly for $\llbracket u^B \rrbracket_{g\mathfrak{E}} = \rho_1$. In the last step of 58, we refer again to eq.29.

$$\begin{aligned} \llbracket z_0 \triangleleft \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} &= \llbracket z_0 \rrbracket_{g\mathfrak{E}} * \llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \\ &= \frac{\left(\llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left(\left(\llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket \lambda^l x_1. (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \right)} \\ &= \frac{(|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{E}} \cdot (|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left((|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{E}} \cdot (|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}} \right)} = |0\rangle_{\mathfrak{E}} \langle 0| \end{aligned} \quad (58)$$

$$\begin{aligned}
& \llbracket y_0 \triangleright (z_0 \triangleleft \lambda^l x_1 \cdot (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2))) \rrbracket_{g\mathfrak{E}} = \llbracket y_0 \rrbracket_{g\mathfrak{E}} * \llbracket (z_0 \triangleleft \lambda^l x_1 \cdot (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2))) \rrbracket_{g\mathfrak{E}} \\
& = \frac{\left(\llbracket z_0 \triangleleft \lambda^l x_1 \cdot (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket z_0 \triangleleft \lambda^l x_1 \cdot (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left(\left(\llbracket z_0 \triangleleft \lambda^l x_1 \cdot (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket z_0 \triangleleft \lambda^l x_1 \cdot (\vee^{\cup} x_1 \triangleright ((x_2 \triangleleft y_2) \triangleright z_2)) \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \right)} \\
& = \frac{(|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{E}} \cdot (|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left((|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{E}} \cdot (|0\rangle_{\mathfrak{E}} \langle 0|)^{\frac{1}{2}} \right)} = |0\rangle_{\mathfrak{E}} \langle 0| \tag{59} \\
& = \llbracket \text{man_die_de_hond_bijt} \rrbracket_{\mathfrak{E}}^1. \tag{60}
\end{aligned}$$

B.2.2 Object Relativization

Just as in the previous derivations, once more we use the interpretations of E_{\setminus} in the two first steps.

$$\llbracket x \triangleright z_2 \rrbracket_{g\mathfrak{E}} = \llbracket x \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} = \frac{\left(\llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket x \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left(\left(\llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket x \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \right)}. \tag{61}$$

$$\begin{aligned}
& \llbracket (x_2 \triangleleft y_2) \triangleright (x \triangleright z_2) \rrbracket_{g\mathfrak{E}} = \llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} * \left(\llbracket x \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right) \\
& = \frac{\left(\llbracket x \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket x \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{E}} \left(\left(\llbracket x \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \cdot \llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} \cdot \left(\llbracket x \rrbracket_{g\mathfrak{E}} * \llbracket z_2 \rrbracket_{g\mathfrak{E}} \right)^{\frac{1}{2}} \right)} \tag{62}
\end{aligned}$$

In the application of the interpretation of *xleft* in eq.4) is the same as in the previous reading, except that controlled comutation is used once, so that $m = 1$, meaning that $(S_+)^1 = S_+$, $U_0 = \mathbb{1}$:

$$T_{\diamond}^0 \left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} \triangleright (x \triangleright z_2) \rrbracket_{g\mathfrak{E}} \right) = \llbracket x \triangleright ((x_2 \triangleleft y_2) \triangleright z_2) \rrbracket_{g\mathfrak{E}} * |0\rangle_{\mathfrak{E}} \langle 0| = |0\rangle_{\mathfrak{E}} \langle 0| \tag{63}$$

$$\begin{aligned}
& \llbracket \lambda^l x_1 \cdot ((x_2 \triangleleft y_2) \triangleright (\vee^{\cup} x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{E}} = \\
& = T_{\square}^0 \left((S_+)^1 T_{\diamond}^0 \left(\llbracket x_2 \triangleleft y_2 \rrbracket_{g\mathfrak{E}} \triangleright (x \triangleright z_2) \rrbracket_{g\mathfrak{E}} \right) \left((S_+)^{\dagger} \right)^1 \right) \\
& = |1\rangle_{\mathfrak{E}} \langle 1| \tag{64}
\end{aligned}$$

Finally, the state ρ_1 is preserved in the interpretations of $E_{/}$ and E_{\setminus} , in the same way that ρ_0 is preserved in the previous reading, according to eq.57, with the result as in eq.30.

$$\begin{aligned}
& \llbracket z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} = \llbracket z_0 \rrbracket_{g\mathfrak{S}} * \llbracket \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \\
& = \frac{\left(\llbracket \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{S}} \cdot \left(\llbracket \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{S}} \left(\left(\llbracket \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{S}} \cdot \left(\llbracket \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}} \right)} \\
& = \frac{(|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{S}} \cdot (|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{S}} \left((|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}} \cdot \llbracket z_0 \rrbracket_{g\mathfrak{S}} \cdot (|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}} \right)} = |1\rangle_{\mathfrak{S}} \langle 1| \tag{65}
\end{aligned}$$

$$\begin{aligned}
& \llbracket y_0 \triangleright (z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2))) \rrbracket_{g\mathfrak{S}} = \llbracket y_0 \rrbracket_{g\mathfrak{S}} * \llbracket z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \\
& = \frac{\left(\llbracket z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{S}} \cdot \left(\llbracket z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{S}} \left(\left(\llbracket z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{S}} \cdot \left(\llbracket z_0 \triangleleft \lambda^l x_1 . ((x_2 \triangleleft y_2) \triangleright (\vee^U x_1 \triangleright z_2)) \rrbracket_{g\mathfrak{S}} \right)^{\frac{1}{2}} \right)} \\
& = \frac{(|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{S}} \cdot (|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}}}{\text{Tr}_{\mathfrak{S}} \left((|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}} \cdot \llbracket y_0 \rrbracket_{g\mathfrak{S}} \cdot (|1\rangle_{\mathfrak{S}} \langle 1|)^{\frac{1}{2}} \right)} = |1\rangle_{\mathfrak{S}} \langle 1| \tag{66}
\end{aligned}$$

$$= \llbracket \text{man_die_de_hond_bijt} \rrbracket_{I\mathfrak{S}}^2. \tag{67}$$

C Connection with Permutation Operator

In the same way that we could go from one interpretation to the other using the covariant and contravariant permutation operators described in [6], we could also assign a permutation operator when we use the positions of words in the sentence as the recipe for contraction. That does not correspond to interpreting directly the derivation steps, but rather to directly following the contraction scheme in figures 3 and 4. To do that, we must assign numbered subsystems to the noun spaces and apply permutation operators in between the partial tracing:

	syn type A	$ A\rangle$
man	n	$[N]^{1*} \otimes \mathfrak{S}$
die	$(n \setminus n) / (\diamond \square np \setminus s)$	$[N]^{1*} \otimes [N]^2 \otimes [S]^* \otimes [N]^3 \otimes \mathfrak{S}$
de hond	np	$[N]^4 \otimes \mathfrak{S}$
bijt	$np \setminus np \setminus s$	$[N]^{4*} \otimes [N]^{3*} \otimes [S] \mathfrak{S}$

Applying partial tracing as direct contraction, we have

$$\begin{aligned}
& \llbracket \text{man_die_de_hond_bijt} \rrbracket_I^1 = \\
& \text{Tr}_1 \left(\llbracket \text{man}^n \rrbracket_I \cdot \text{Tr}_3 \left(\llbracket \text{die}^{(n \setminus n) / (\diamond \square np \setminus s)} \rrbracket_I \cdot \text{Tr}_4 \left(\llbracket \text{de_hond}^{np} \rrbracket_I \cdot \llbracket \text{bijt}^{np \setminus np \setminus s} \rrbracket_I \right) \right) \right). \tag{68}
\end{aligned}$$

Applying permutation operators to obtain the interpretation of the second reading, we have

$$\begin{aligned}
& \llbracket \text{man_die_de_hond_bijt} \rrbracket_I^2 = \\
& \text{Tr}_1 \left(\llbracket \text{man}^n \rrbracket_I \cdot \text{Tr}_3 \left(\llbracket \text{die}^{(n \setminus n) / (\diamond \square np \setminus s)} \rrbracket_I \cdot \text{Tr}_4 \left(\llbracket \text{de_hond}^{np} \rrbracket_I \cdot P^{34} \llbracket \text{bijt}^{np \setminus np \setminus s} \rrbracket_I P^{34} \right) \right) \right). \tag{69}
\end{aligned}$$

D Proof transformation: β reduction

The β -reduction is one of the rewrite rules of the λ -calculus. It asserts that applying a term with a lambda-bound variable to a certain argument is equivalent to substituting that argument directly in the original term, before introducing the lambda. In proof-theoretic terms, if an introduction rule is used followed by an elimination rule, the derivation is not minimal. To elucidate this point, below is the skeleton of a derivation where a term of type A is proved twice, by axiom and by an unknown proof:

$$\frac{\frac{\frac{\overline{x : A \vdash x : A} \text{ axiom}}{\vdots}}{\Delta \vdash n : A} \quad \frac{\frac{\overline{x : A, \Gamma \vdash t : B}}{\Gamma \vdash \lambda^l x.m : A \setminus B} \setminus I}{(\Gamma, \Delta) \vdash n \triangleright (\lambda^l x.m) : B} \setminus E}{\Delta, \Gamma \vdash m[x/n] : B} .$$

The β reduction consists of substituting the unknown proof of the term of type A in place of the axiom, reducing the need for the double proof of that term, and consequently the size of the proof:

$$\frac{\frac{\vdots}{\Delta \vdash n : A}}{\Delta, \Gamma \vdash m[x/n] : B} .$$

Through this reduction, a map from one conclusion to the other can be obtained, which has to be an equality regarding their interpretations:

$$\llbracket n \triangleright (\lambda^l x.m) \rrbracket_g = \llbracket m[x/n] \rrbracket_g, \forall g.$$

This equality will be used to check that the density matrix construction interpretation is consistent with the λ -calculus. Below a concrete symbolic derivation before the reduction is shown:

$$\frac{\frac{\frac{\overline{w : B \vdash w : B} \text{ ax}}{w : B, z : B \setminus (A/B) \vdash (w \triangleright z) : A/B} \setminus E_2 \quad \frac{\frac{\overline{x : A/B \vdash x : A/B} \text{ ax} \quad \overline{y : B \vdash y : B} \text{ ax}}{x : A/B, y : B \vdash (x \triangleleft y) : A} \setminus E_1}{y : B \vdash \lambda^l x.(x \triangleleft y) : (A/B) \setminus A} \setminus I_1}{(w : B, z : B \setminus (A/B), u : B) \vdash (w \triangleright z) \triangleright (\lambda^l x.(x \triangleleft y)) : A} \setminus E_3} .$$

The interpretation in the spatial space S of the several steps of the proof is given below, following the numbering in the proof:

$$E_{/1} : \llbracket (x \triangleleft y) \rrbracket_{g^S} = \sum_{ii', jj'} S \mathbf{X}_{jj'}^{ii} S \mathbf{Y}^{jj} |i\rangle_{[A]} \langle i'|,$$

$$I_{\setminus 1} : \llbracket \lambda^l x.(x \triangleleft y) \rrbracket_{g^S} = \sum_{ii', jj'} |i'\rangle_{[B]} \langle i|_{[A]^*} \otimes S \mathbf{Y}^{jj} |i\rangle_{[A]} \langle i'|,$$

$$E_{\setminus 2} : \llbracket (w \triangleright z) \rrbracket_g = \sum_{ll', mm', nn'} S \mathbf{W}^{ll'} S \mathbf{Z}_{ll', nn'}^{mm'} |m'\rangle_{[A] \otimes [B]^*} \langle n|,$$

$$E_{\setminus 3} : \llbracket (w \triangleright z) \triangleright (\lambda^l x. (x \triangleleft y)) \rrbracket_{g^S} = \sum_{ii', jj', ll'} S \mathbf{W}^{ll'} S \mathbf{Z}^{ii'} S \mathbf{Y}^{jj'} |i\rangle_{[A]} \langle i'|.$$

In spin space \mathfrak{S} the interpretation of the proof steps is as follows:

$$E_{/1} : \llbracket (x \triangleleft y) \rrbracket_{g^{\mathfrak{S}}} = \llbracket x \rrbracket_{g^{\mathfrak{S}}} * \llbracket y \rrbracket_{g^{\mathfrak{S}}}$$

$$I_{\setminus 1} : \llbracket \lambda^l x. (x \triangleleft y) \rrbracket_{g^{\mathfrak{S}}} = I * \llbracket y \rrbracket_{g^{\mathfrak{S}}} = \llbracket y \rrbracket_{g^{\mathfrak{S}}}$$

$$E_{\setminus 2} : \llbracket (w \triangleright z) \rrbracket_g = \llbracket w \rrbracket_{g^{\mathfrak{S}}} * \llbracket z \rrbracket_{g^{\mathfrak{S}}}$$

$$E_{\setminus 3} : \llbracket (w \triangleright z) \triangleright (\lambda^l x. (x \triangleleft y)) \rrbracket_{g^S} = \left(\llbracket w \rrbracket_{g^{\mathfrak{S}}} * \llbracket z \rrbracket_{g^{\mathfrak{S}}} \right) * \llbracket y \rrbracket_{g^{\mathfrak{S}}}$$

A similar treatment is done for the derivation after the reduction:

$$\frac{\frac{w : B \vdash w : B}{ax} \quad \frac{z : B \setminus (A/B) \vdash z : B \setminus (A/B)}{ax}}{\frac{w : B, z : B \setminus (A/B) \vdash (w \triangleright z) : A/B}{\setminus E_2} \quad \frac{y : B \vdash y : B}{ax}}{\frac{w : B, z : B \setminus (A/B), u : B \vdash ((w \triangleright z) \triangleleft y) : A}{\setminus E_4}}.$$

The value of $\llbracket (w \triangleright z) \rrbracket_g$ is the same as before. For $\llbracket ((w \triangleright z) \triangleleft y) \rrbracket_{g^S}$:

$$E_{\setminus 4} : \llbracket ((w \triangleright z) \triangleleft y) \rrbracket_{g^S} = \sum_{ii', jj', ll'} S \mathbf{W}^{ll'} S \mathbf{Z}^{ii'} S \mathbf{Y}^{jj'} |i\rangle_{[A]} \langle i'|.$$

On the spin space, we have

$$E_{\setminus 4} : \llbracket ((w \triangleright z) \triangleleft y) \rrbracket_{g^{\mathfrak{S}}} = \left(\llbracket w \rrbracket_{g^{\mathfrak{S}}} * \llbracket z \rrbracket_{g^{\mathfrak{S}}} \right) * \llbracket y \rrbracket_{g^{\mathfrak{S}}}$$

Comparing the two derivations and interpretations, the conclusion is that

$$\llbracket E_{\setminus 4}(y, z(w)) \rrbracket_{g^S} = \llbracket E_{\setminus 3}(z(w), \lambda x. x(y)) \rrbracket_{g^S},$$

as expected, and

$$\llbracket E_{\setminus 4}(y, z(w)) \rrbracket_{g^{\mathfrak{S}}} = \llbracket E_{\setminus 3}(z(w), \lambda x. x(y)) \rrbracket_{g^{\mathfrak{S}}}.$$